The Geometry of the Dot and Cross Products

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1 Introduction

Most students first learn the algebraic formula for the dot and cross products in rectangular coordinates, and only then are shown their geometric interpretations. We believe this should be done in the other order. Students tend to remember best the first definition they use; this should not be an algebraic formula devoid of context. The geometric definition is coordinate independent, and therefore conveys invariant properties of these products, not just a formula for calculating them. Furthermore, it is easier to derive the algebraic formula from the geometric one than the other way around, as we demonstrate below.

2 Dot Product

The dot product is fundamentally a projection. As shown in Figure 1, the dot product of a vector with a unit vector is the projection of that vector in the direction given by the unit vector. This leads to the geometric formula

$$\vec{v} \cdot \vec{w} = |\vec{v}| |\vec{w}| \cos \theta$$

for the dot product of any two vectors $\vec{v}$ and $\vec{w}$.

An immediate consequence of (1) is that the dot product of a vector with itself gives the square of the length, that is

$$\vec{v} \cdot \vec{v} = |\vec{v}|^2$$
In particular, taking the “square” of any unit vector yields 1, for example
\[ \hat{i} \cdot \hat{i} = 1 \] (3)
where \( \hat{i} \) as usual denotes the unit vector in the \( x \) direction. \(^{1}\) Furthermore, it follows immediately from the geometric definition that two vectors are orthogonal if and only if their dot product vanishes, that is
\[ \vec{v} \perp \vec{w} \iff \vec{v} \cdot \vec{w} = 0 \] (4)

For instance, if \( \hat{j} \) denotes the unit vector in the \( y \) direction, then
\[ \hat{i} \cdot \hat{j} = 0 \] (5)

The geometry of an orthonormal basis is fully captured by these properties; each basis vector is normalized, which is (3), and each pair of vectors is orthogonal, which is (5).

The components of a vector \( \vec{v} \) in an orthonormal basis are just the dot products of \( \vec{v} \) with each basis vector. For instance, in two dimensions, setting
\[ v_x = \vec{v} \cdot \hat{i} \]
\[ v_y = \vec{v} \cdot \hat{j} \] (6)

\(^{1}\)We follow standard usage among scientists and engineers by putting hats on unit vectors.
implies $\vec{v} = v_x \hat{i} + v_y \hat{j}$. The component form of the dot product now follows from its properties given above. For example, if $\vec{w} = w_x \hat{i} + w_y \hat{j}$, then

$$
\vec{v} \cdot \vec{w} = (v_x \hat{i} + v_y \hat{j}) \cdot (w_x \hat{i} + w_y \hat{j}) \\
= v_x w_x \hat{i} \cdot \hat{i} + v_y w_y \hat{j} \cdot \hat{j} + v_x w_y \hat{i} \cdot \hat{j} + v_y w_x \hat{j} \cdot \hat{i} \\
= v_x w_x + v_y w_y
$$

This computation clearly works for any orthonormal basis. A special case is the dot product of a vector with itself, which reduces to the Pythagorean theorem, for example

$$
\vec{v} \cdot \vec{v} = |\vec{v}|^2 = v_x^2 + v_y^2 \tag{8}
$$

What happens if you don’t use an orthonormal basis? Consider Figure 2, in which $\vec{A} + \vec{C} = \vec{B}$, or equivalently $\vec{C} = \vec{B} - \vec{A}$. Then

$$
\vec{C} \cdot \vec{C} = (-\vec{A} + \vec{B}) \cdot (-\vec{A} + \vec{B}) \\
= \vec{A} \cdot \vec{A} + \vec{B} \cdot \vec{B} - 2 \vec{A} \cdot \vec{B} \tag{9}
$$

or equivalently

$$
|\vec{C}|^2 = |\vec{A}|^2 + |\vec{B}|^2 - 2|\vec{A}||\vec{B}| \cos \theta \tag{10}
$$

which is just the Law of Cosines! The Law of Cosines is usually used to derive the geometric form of the dot product (1) from the algebraic form (7), which is taken as the definition. Instead, by starting with geometry, the Law of Cosines follows immediately.
Not so fast! Did you spot the flaw in the above argument? In the computation (7) of the algebraic formula for the dot product in terms of components, it was assumed without comment that the dot product distributes over addition, or in other words that the dot product is linear. If one starts with the geometric definition (1), this must be proved.

However, the proof is straightforward, as shown in Figure 3. We must show that
\[(\vec{v} + \vec{u}) \cdot \vec{w} = \vec{v} \cdot \vec{w} + \vec{u} \cdot \vec{w}\]  
(11)

But this is equivalent to showing that the projection of $\vec{v} + \vec{u}$ along $\vec{w}$ is the sum of the projections of $\vec{v}$ and $\vec{u}$, which is immediately obvious from Figure 3.

3 Examples

What is the bonding angle of carbon tetrachloride? Take a tetrahedron, and connect each vertex (a chlorine atom) to the center (a carbon atom). What is the angle between the lines that meet at the center?

This problem can be done by brute force using high school geometry — try it. A simpler approach is to represent the tetrahedron using vectors. It helps to realize that a tetrahedron is formed by connecting alternating vertices of a cube, as shown in Figure 4, and that the center of the tetrahedron is at the center of the cube. It is now straightforward to write down the coordinates of the vertices, thus obtaining the components of the vectors from the center.

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$^2$Active versions of this figure are available online at [1] in both Java and Maple formats.
(shown in red in the figure). Compute the dot product (algebraically), divide by the magnitudes, and use the geometric expression for the dot product to read off the cosine of the angle. The answer? A bit more than 109°.

A simpler version of this problem would be to determine the angle between the diagonal of a cube and one of its edges, that is, between the green (or red) lines and the adjacent black lines in Figure 4. Or between the diagonals of adjacent faces — the blue lines in the figure. Again, both the geometric and algebraic expressions for the dot product are involved in the solution.

4 Gram-Schmidt Orthogonalization

The most common use of the dot product in applications in physics and engineering is to decompose vectors into their components parallel and perpendicular to a given vector, for which an understanding of the geometric definition (1) is essential.

The Gram-Schmidt orthogonalization process uses this idea to construct an orthonormal basis from a given set of (linearly independent) vectors. This can be beautifully illustrated in three dimensions, as shown in Figure 5 and described below. Stick out three fingers of one hand in arbitrary directions. Pick one of these vectors (black) as the first of the new orthogonal basis vectors. Then pick a second (thick red), and subtract from this vector its projection parallel to the first, resulting in a vector perpendicular to the first (thin red). Now subtract from the remaining vector (thick blue) its projections parallel to both the first and second vectors, resulting in a vector perpendicular to both (thinnest blue). If your fingers are limber enough, you
have an orthogonal basis. These three vectors (black, thin red, thinnest blue) can be normalized if desired, yielding an orthonormal basis. Students who see this geometric idea behind the Gram-Schmidt process have a much easier time sorting their way through the morass of formulas in textbooks.

In advanced courses, the fact that two vectors are perpendicular if their dot product is zero may be used in more abstract settings, such as Fourier analysis. A problem which asks students to find the vector perpendicular to a given vector, first in two and then in three dimensions, provides an excellent introduction to this idea.

5 Cross Product

The cross product is fundamentally a directed area. The magnitude of the cross product is defined to be the area of the parallelogram shown in Figure 6. This leads to the formula

\[ |\vec{v} \times \vec{w}| = |\vec{v}||\vec{w}| \sin \theta \]

an immediate consequence of which is that

\[ \vec{v} \parallel \vec{w} \iff \vec{v} \times \vec{w} = \vec{0} \]

The direction of the cross product is given by the right-hand rule, so that in the example shown \( \vec{v} \times \vec{w} \) points into the page. This implies that

\[ \vec{v} \times \vec{w} = -\vec{w} \times \vec{v} \]
so that the cross product is not commutative. 3 Another important property of the cross product is that
\[ \vec{v} \times \vec{v} = \vec{0} \]
which also follows immediately from (12).

In terms of the standard orthonormal basis, the geometric formula quickly yields
\[
\begin{align*}
\hat{i} \times \hat{j} &= \hat{k} \\
\hat{j} \times \hat{k} &= \hat{i} \\
\hat{k} \times \hat{i} &= \hat{j}
\end{align*}
\]
(16)
with the remaining products being determined by (14) and (15). This cyclic nature of the cross product can be emphasized by diagramming the multiplication table as shown in Figure 7. 4 Products in the direction of the arrow get a plus sign (e.g. \( \hat{k} \cdot \hat{i} = +\hat{j} \)), while products against the arrow get a minus sign (e.g. \( \hat{i} \cdot \hat{k} = -\hat{j} \)).

3This may be the first example some students have seen of a product which is not commutative. It is worth pointing out that the cross product also fails to be associative, since for example \( \hat{i} \times (\hat{i} \times \hat{j}) = -\hat{j} \) but \( (\hat{i} \times \hat{i}) \times \hat{j} = \vec{0} \).

4This is really the multiplication table for the unit imaginary quaternions, a number system which generalizes the familiar complex numbers. Quaternions predate vector analysis, which borrowed the quaternionic units \( i, j, k \) as labels for the rectangular basis vectors \( \hat{x}, \hat{y}, \hat{z} \). A more logical name for the rectangular basis vectors would be \( \hat{x}, \hat{y}, \hat{z} \), which is used by many physicists.
Using an orthonormal basis such as \( \{ \hat{i}, \hat{j}, \hat{k} \} \), the geometric formula reduces to the standard component form of the cross product. If \( \vec{v} = v_x \hat{i} + v_y \hat{j} + v_z \hat{k} \) and and \( \vec{w} = w_x \hat{i} + w_y \hat{j} + w_z \hat{k} \), then

\[
\vec{v} \times \vec{w} = (v_x \hat{i} + v_y \hat{j} + v_z \hat{k}) \times (w_x \hat{i} + w_y \hat{j} + w_z \hat{k})
\]

\[
= v_x w_x \hat{i} \times \hat{i} + v_x w_y \hat{i} \times \hat{j} + ... 
\]

\[
= (v_y w_z - v_z w_y) \hat{i} + (v_z w_x - v_x w_z) \hat{j} + (v_x w_y - v_y w_x) \hat{k}
\]

which is often written as the symbolic determinant

\[
\vec{v} \times \vec{w} = \begin{vmatrix} \hat{i} & \hat{j} & \hat{k} \\ v_x & v_y & v_z \\ w_x & w_y & w_z \end{vmatrix}
\]

(18)

As with the dot product, this derivation works in any (right-handed) orthonormal basis.

OK, this time you surely noticed; we must again check linearity, that is, whether the cross product distributes over addition, or in other words whether

\[
\vec{w} \times (\vec{v} + \vec{u}) = \vec{w} \times \vec{v} + \vec{w} \times \vec{u}
\]

(19)

As we now show, this follows with a little thought from Figure 8. Consider in turn the vectors \( \vec{v}, \vec{u}, \) and \( \vec{v} + \vec{u} \). The cross product of each of these vectors with \( \vec{w} \) is proportional to its projection perpendicular to \( \vec{w} \). These projections are shown as solid lines in the figure. Since the projections lie in the plane perpendicular to \( \vec{w} \), they can be combined into the triangle shown in the middle of the figure. Two of the vectors making up the sides of a
Figure 8: A geometric proof of the linearity of the cross product.

triangle add up to the third; in this case, the sides are the projections of \( \vec{v} \), \( \vec{u} \), and \( \vec{v} + \vec{u} \), and the latter is clearly the sum of the first two. But each cross product is now just a rotation of one of the sides of this triangle, rescaled by the length of \( \vec{w} \); these are the arrows perpendicular to the faces of the prism. Two of these vectors therefore still add to the third, as indicated by the vector triangle in front of the prism. This establishes (19).

6 Cross Product Pedagogy

We digress briefly to mention some pedagogical issues which arise when teaching students about the cross product.

First of all, we strongly discourage teaching, or even reviewing, the dot and cross products at the same time — students tend to get them mixed up!

As for the calculation of the cross product, we encourage students to compute the determinant (18), rather than memorizing (17). But we also encourage students to use the multiplication table directly for simple cross products, such as \((\hat{i} + 3\hat{j}) \times \hat{k}\), rather than using determinants at all.

However, we discourage students from working out the determinant using minors, which would result in the final expression in (17) but with two sign changes in the term involving \( \hat{j} \). Instead, we encourage students to take \( 3 \times 3 \)
determinants in the form
\[
\begin{vmatrix}
\hat{i} & \hat{j} & \hat{k} \\
v_x & v_y & v_z \\
w_x & w_y & w_z
\end{vmatrix}
\]

(20)

where one multiplies the terms along each diagonal line, subtracting the products obtained along lines going down to the left from those along lines going down to the right. While this method works only for (2 × 2 and) 3 × 3 determinants, and is therefore usually omitted from a linear algebra course, it has the definite advantage of emphasizing the cyclic nature of the cross product — whereas students who use minors often make sign errors.

7 Divergence and Curl

The geometric approach extends naturally to the divergence and curl of a vector field.

It is straightforward to define divergence as flux per unit volume. This amounts to taking an infinitesimal version of the Divergence Theorem as the definition of divergence; the computation traditionally done to prove the theorem become instead the derivation of the component representation of the divergence, expressed in rectangular coordinates.

This approach is beautifully summarized in [2], and is included in some physics textbooks, such as [3, 4, 5], but has not been adopted in many mathematics textbooks (a notable exception being [6]). It not only emphasizes the geometric relationship between divergence and flux, but makes clear how to derive the component representation in other coordinate systems.

It then seems natural to define curl as (oriented) circulation per unit area, turning an infinitesimal version of Stokes’ Theorem into a definition, which is then used to derive the component representation of curl in rectangular coordinates. This is indeed the approach taken in [2] (and [6]).

However, there is an issue of principle which we have not seen addressed in textbooks: One must argue that this procedure defines a vector field! It is not enough to derive the separate formulas for the circulation per unit area in each (rectangular) coordinate plane, one must also show that the circulation in any plane is equal to an appropriate sum of these circulations. But this follows immediately from the diagram in Figure 9, which decomposes an arbitrary triangle into three triangles in the coordinate planes; it
is straightforward to check that the circulation around the original triangle is equal to the sum of the (appropriately oriented!) circulations around the “component” triangles. The contributions of the circulations along all the “extra” sides cancel pairwise.

8 Discussion

We have shown how to prove geometrically that the dot and cross products are linear (i.e. that they distribute over addition), using Figures 3 and 8, respectively. These proofs are important for faculty wishing to emphasize the geometric definitions in their classroom. However, it is not at all clear that these proofs are relevant for students, especially in introductory courses. Most students will accept, say, the argument presented in (7) without question, failing to notice that linearity is an issue. While faculty should of course be prepared to justify this if asked, we see nothing wrong with not raising the issue of linearity otherwise. Similarly, the proof that the curl is a vector field can be omitted.

As shown above, the Law of Cosines follows immediately from the geometric definition of the dot product, in direct contrast to the traditional treatment, in which the order is reversed. Indeed, many students who have memorized the Law of Cosines are delighted with this derivation. This suggests to us that it might be beneficial to include the dot product in trigonometry courses, even in high school, rather than saving it for later. A strong argument in favor of such an approach is the ease with which the addition for-
formulas could then be derived: From the unit circle definition of the trig functions, any vector on the unit circle can be written in the form $\cos \theta \hat{i} + \sin \theta \hat{j}$, where $\theta$ is the angle to the positive $x$-axis. But the cosine of the angle between any two such vectors is just their dot product, since both are unit vectors, so that

$$
\cos(\beta - \alpha) = \vec{u} \cdot \vec{v} = (\cos \alpha \hat{i} + \sin \alpha \hat{j}) \cdot (\cos \beta \hat{i} + \sin \beta \hat{j})
$$

$$
= \cos \alpha \cos \beta + \sin \alpha \sin \beta
$$

(21)

as illustrated in Figure 10.

Why aren’t geometric proofs such as those presented here more common? We believe this lies in part with the difficulty in translating them into words on the printed page. Consider for example the Gram-Schmidt orthogonalization process. As discussed in Section 4, the geometry of this construction can be beautifully illustrated by (literally) waving one’s hands. It is easy to write down the formulas after seeing this demonstration, but the formulas by themselves are much less enlightening. You may have had some trouble sorting through the arguments in this paper, but try them in the classroom, where you can point at the pictures. They work like a charm.

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References


